

LiMnVO₄

Masahiko Sugahara,* Akira
Yoshiasa, Takamitsu Yamanaka
and Humihiko Takei

Department of Earth and Space Science,
Graduate School of Science, Osaka University,
1-1 Machikaneyama, Toyonaka, Osaka
560-0043, Japan

Correspondence e-mail:
masa-s@ess.sci.osaka-u.ac.jp

Key indicators

Single-crystal X-ray study
T = 297 K
Mean $\sigma(V-O)$ = 0.001 Å
R factor = 0.022
wR factor = 0.033
Data-to-parameter ratio = 64.0

For details of how these key indicators were
automatically derived from the article, see
<http://journals.iucr.org/e>.

The crystal structure of lithium manganese vanadate, LiMnVO₄, is characterized by chains of edge-sharing MnO₆ octahedra, bridged by pairs of edge-sharing LiO₄ and VO₄ tetrahedra. Cation–cation repulsions displace the V and Li ions from the tetrahedral centres and deform the tetrahedra further by shortening the mutually shared O–O edges. The smallest r.m.s. atomic displacements in thermal motion for Li and V are parallel to the shortest Li···V vector with maximum vibrational motion normal to that direction. Cation–cation interactions between the Li and V tetrahedral sites play an important role.

Comment

Li_xMn₂O₄ (0 < x < 2) and vanadium oxides such as V₂O₅ and LiV₃O₈ have been widely investigated as candidates for cathode materials in lithium batteries. The LiMnVO₄ phase has been found in the pseudo-binary system of LiMn₂O₄–LiV₂O₄. The structure of the phase was proposed by Rietveld analysis (Sato & Kano, 1994), and consists of an arrangement of edge-sharing MnO₆ octahedral chains, linked by VO₄ and LiO₄ tetrahedra sharing a common edge.

The lithium ion, which has a low mass and is monovalent, usually exhibits a large anisotropic thermal motion in ionic conductors. However, although the displacement modes contain much information about possible lithium migration

Received 23 September 2003

Accepted 20 October 2003

Online 8 November 2003

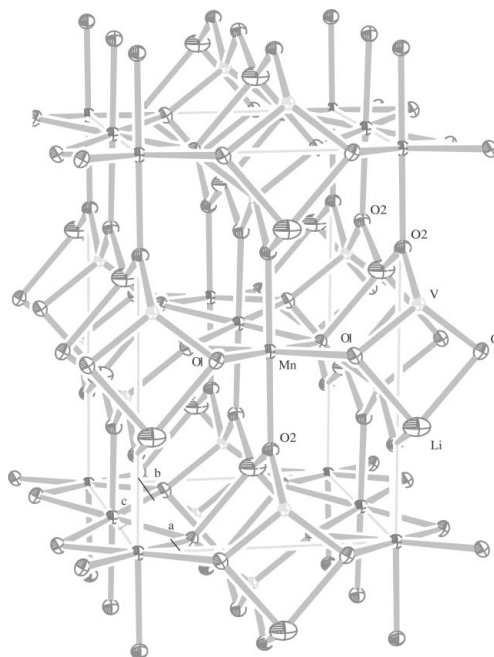


Figure 1

Perspective view of the LiMnVO₄ structure. Displacement ellipsoids are drawn at the 50% probability level.

paths, no detailed single-crystal structure analysis of LiMnVO₄ has been published. In this study, we have succeeded in synthesizing a large single crystal of LiMnVO₄ and the structure refinement with full anisotropic displacement parameters is reported. The large anisotropic thermal motion of Li in the edge-sharing LiO₄ tetrahedral site is discussed.

Sato & Kano (1994), Sato *et al.* (1996), and Padhi *et al.* (1997) have investigated the crystal structure of LiMnVO₄ via the Rietveld method, with an isotropic displacement model for all atoms. We find good agreement between the atomic coordinates of their powder determination and those of our single-crystal study.

The crystal structure of LiMnVO₄ (Fig. 1) is characterized by chains of edge-sharing MnO₆ octahedra, running parallel to the *c* axis. These chains are bridged by pairs of LiO₄ and VO₄ tetrahedra. Both tetrahedral sites are linked to each other by sharing an edge. Both the Mn···Mn [3.1815 (2) Å] and V···Li [2.684 (6) Å] distances parallel to the *c* and *b* axes, respectively, are relatively short, resulting in a strong repulsion between cations.

Interatomic distances for the MnO₆, VO₄ and LiO₄ polyhedra are listed in Table 1. The distances of O1···O1^{iv} between MnO₆ octahedra themselves and O1···O1ⁱⁱⁱ shared-edges between VO₄ and LiO₄ tetrahedra are 2.992 (1) and 2.772 (1) Å, respectively (symmetry codes as in Table 1). The O1···O1ⁱⁱⁱ edges are shortened, due to cation–cation repulsion, which accompanies the significant off-centering of V and Li in each tetrahedral site.

An additional measure of the distortion of the VO₄ and LiO₄ tetrahedra is given by the deviation of the bond angles from the ideal 109.48°.

Fig. 1 shows the r.m.s. displacements of Mn, V, Li and O atoms rendered as ellipsoids. The smallest Li atom mean atomic displacements are oriented along the Li···V vectors and the maximal displacement vector is normal to this direction. The cation–cation interaction between the Li and V tetrahedral sites play an important role.

Many Li cathode materials, such as the spinel-type structure and layer structure (Sato & Kano, 1994), have an Li-ion conduction path in the structure. On the other hand, the Li sites in this LiMnVO₄ structure are widely separated from each other and no vacant site which the Li ion could occupy is located near the Li site, though the amplitude and anisotropy of thermal vibration for Li is large enough. It is, therefore, inferred that this phase does not suit a cathode material. This phase transforms to a spinel-type phase (Padhi *et al.*, 1997) under high pressure and temperature. The spinel-type phase should be appropriate for a cathode material. It is emphasized that the largest amplitude of the Li ion in this phase corresponds to the direction of displacement in the transition from this phase to the spinel-type phase, and the transition is achieved only by the displacement of Li ions.

Experimental

A single crystal of LiMnVO₄ was grown using an LiVO₃ flux. A mixture of Li₂CO₃ (0.1 mol), V₂O₅ (0.1 mol) and MnCO₃ (0.05 mol)

was heated at 1570 K for 24 h under atmospheric conditions. After slow cooling to 1070 K at a rate of 1 K h⁻¹, the product was quenched. A flat LiMnVO₄ single crystal of 10 × 10 × 1 mm was ground into a sphere. The stoichiometric composition was confirmed by an ICP emission spectrochemical analysis.

Crystal data

| | |
|---|-------------------------------------|
| LiMnVO ₄ | Mo Kα radiation |
| <i>M_r</i> = 176.82 | Cell parameters from 25 reflections |
| Orthorhombic, <i>Cmcm</i> | $\theta = 39.4\text{--}49.0^\circ$ |
| <i>a</i> = 5.7640 (3) Å | $\mu = 6.68\text{ mm}^{-1}$ |
| <i>b</i> = 8.7418 (9) Å | <i>T</i> = 297 K |
| <i>c</i> = 6.3629 (3) Å | Sphere, brown |
| <i>V</i> = 320.61 (4) Å ³ | 0.08 mm (radius) |
| <i>Z</i> = 4 | |
| <i>D_x</i> = 3.663 Mg m ⁻³ | |

Data collection

| | |
|--|--|
| Rigaku AFC-5S diffractometer | <i>R</i> _{int} = 0.014 |
| $\theta/2\theta$ scans | $\theta_{\text{max}} = 50.0^\circ$ |
| Absorption correction: sphere (International Tables, Vol C, Table 6.3.3.3) | <i>h</i> = 0 → 12 |
| <i>T</i> _{min} = 0.330, <i>T</i> _{max} = 0.343 | <i>k</i> = -18 → 0 |
| 1759 measured reflections | <i>l</i> = -13 → 13 |
| 944 independent reflections | 3 standard reflections every 150 reflections |
| 902 reflections with <i>F</i> > 3σ(<i>F</i>) | intensity decay: none |

Refinement

| | |
|---|---|
| Refinement on <i>F</i> | (Δ/σ) _{max} < 0.001 |
| <i>R</i> = 0.022 | $\Delta\rho_{\text{max}} = 0.18\text{ e \AA}^{-3}$ |
| <i>wR</i> = 0.033 | $\Delta\rho_{\text{min}} = -1.19\text{ e \AA}^{-3}$ |
| <i>S</i> = 1.95 | Extinction correction: RADY89 (Sasaki, 1989) |
| 1663 reflections | Extinction coefficient: 0.057 (4) |
| 26 parameters | |
| $w = 1/[\sigma^2(F_o) + (0.013936F_o)^2]$ | |

Table 1

Selected geometric parameters (Å, °).

| | | | |
|--------------------------------------|------------|------------------------|-------------|
| Li–O1 | 2.123 (4) | Mn–O2 | 2.1602 (10) |
| Li–O2 ⁱ | 1.962 (2) | V–O1 | 1.7544 (9) |
| Mn–O1 ⁱⁱ | 2.1909 (6) | V–O2 | 1.6896 (10) |
| O1–Li–O1 ⁱⁱⁱ | 81.5 (2) | O1 ^v –Mn–O2 | 90.65 (6) |
| O1–Li–O2 ^{iv} | 107.7 (2) | O1–V–O1 ⁱⁱⁱ | 104.35 (10) |
| O2 ⁱ –Li–O2 ^{iv} | 132.5 (3) | O1–V–O2 | 110.55 (2) |
| O1 ^v –Mn–O1 ^{vi} | 86.12 (6) | O2–V–O2 ^{vii} | 110.2 (1) |

Symmetry codes: (i) $-x, 1 - y, \frac{1}{2} + z$; (ii) $\frac{1}{2} - x, \frac{1}{2} - y, -z$; (iii) $-x, y, \frac{1}{2} - z$; (iv) $x, 1 - y, -z$; (v) $\frac{1}{2} - x, y - \frac{1}{2}, \frac{1}{2} - z$; (vi) $x - \frac{1}{2}, y - \frac{1}{2}, z$; (vii) $x, y, \frac{1}{2} - z$.

Precession photographs showed an orthorhombic unit cell with systematic absences of *hkl*; *h* + *k* = 2*n* + 1 and *h0l*; *h* and *l* = 2*n* + 1. The possible space groups were therefore *Cmcm*, *Cmc2₁* and *C2cm*. The space group was initially assumed to be *Cmcm* before the remaining two space groups were tested. Calculations gave the reliability factor *R* = 0.0217 for *Cmcm*, *R* = 0.0217 for *Cmc2₁* and *R* = 0.0223 for *C2cm*, all with anisotropic displacement parameters. Since all structure models converged to the *Cmcm* model and were essentially identical, within errors, the space group *Cmcm* was adopted. The minimum electron-density peak was located near the position (0, 0.08, 0). All atoms are located on special positions. Crystallographically equivalent reflections were averaged.

Data collection: *MSC Software* (Rigaku, 1996); cell refinement: *MSC Software*; data reduction: *MSC Software*; program(s) used to solve structure: *RADY89* (Sasaki, 1989); program(s) used to refine structure: *RADY89*; molecular graphics: *ORTEP-3 for Windows* (Farrugia, 1997).

We thank Associate Professor Nobuo Ishizawa for his help.

References

Farrugia, L. J. (1997). *J. Appl. Cryst.* **30**, 565.

- Padhi, A. K., Archibald, W. B., Nanjundaswamy, K. S. & Goodenough, J. B. (1997). *J. Solid State Chem.* **128**, 267–272.
- Rigaku (1996). *MSC Software*. Version 4.4.1. Rigaku, Tokyo, Japan.
- Sasaki, S. (1989). *RADY89*. Tokyo Institute of Technology, Japan.
- Sato, M. & Kano, S. (1994). *Chem. Lett.* **3**, 427–430.
- Sato, M., Kano, S., Tamaki, S., Misawa, M., Shirakawa, Y. & Ohashi, M. (1996). *J. Mater. Chem.* **6**, 1191–1194.